

# Steep Sub-Boltzmann Switching in AlGaIn/GaN Phase-FETs With ALD VO<sub>2</sub>

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**Abstract**—We report the first demonstration of ultralow leakage AlGaIn/GaN MOS-HEMTs on silicon substrates loaded at the source with atomic layer deposition (ALD)-grown VO<sub>2</sub> metal-insulator transition (MIT) material resistors. The resulting GaN phase-transition FET (phase-FET) shows steep sub-Boltzmann switching in its transfer characteristics in both sweep directions with slopes of  $\sim 9.9$  and  $\sim 28.2$  mV/decade at 60 °C. The control MOS-HEMTs and the phase-FETs show more than 12-order on/off ratio with ultralow off-state leakage. Because of the MIT, the drain current and the transconductance also show unconventional behavior. This first demonstration of ultralow leakage steep switching in GaN phase-FETs using integration-friendly ALD VO<sub>2</sub> opens the door to introducing new functionalities in nitride low-power digital devices, microwave circuits, photonic devices, and power electronics in the GaN-on-silicon platform.

**Index Terms**—Atomic layer deposition (ALD), Boltzmann limit, GaN, hyper-FETs, phase-transition FET (phase-FETs), steep switching.

## I. INTRODUCTION

TO CONTINUE the march of Moore's law, an increasing number of non-Si materials such as high- $k$  gate dielectrics and strained SiGe are now used in traditional CMOS technology [1], [2]. The III-V semiconductor heterostructures and 2-D materials are being explored for the next-generation devices and circuits [3]–[6]. Transition-metal oxide materials are also being investigated because they host many forms of electronic, structural, and magnetic phase transitions. Such phase transitions offer the potential to realize

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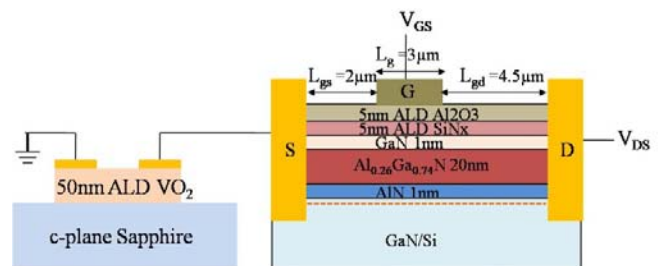


Fig. 1. Schematic of realized phase-FET. AlGaIn/GaN MOS-HEMT on Si (right) loaded at the source with ALD VO<sub>2</sub> resistor (left:  $2 \mu\text{m} \times 100 \mu\text{m}$ ).

new device architectures and can help to overcome some of the fundamental limitations of traditional transistors, such as the  $\sim 60$  mV/decade Boltzmann limit of subthreshold switching at room temperature [7], [8]. Steep switching is currently being investigated by combining ferroelectric materials in the gate-stack of FETs to realize negative-capacitance FETs [9]. Another method is by loading the source electrode of Boltzmann FETs with materials exhibiting reversible metal-insulator transitions (MITs) to realize phase-transition FETs (phase-FETs/hyper-FETs) [10], [11]. Phase-FETs with steep switching and enhanced ON currents were proposed and demonstrated in 2015 by combining molecular beam epitaxy grown VO<sub>2</sub> with Si FETs [10].

VO<sub>2</sub> is an oxide material that undergoes a reversible temperature-driven structural and electronic MIT at  $\sim 67$  °C accompanied with up to five orders of magnitude change in resistivity [12]. This phase transition can also be induced electronically through an applied bias, or optically by excitation from a photonic source [10], [13]. Recently, we demonstrated steep switching in GaN FETs by combining atomic layer deposition (ALD) VO<sub>2</sub> with AlGaIn/GaN HEMTs [14]. ALD VO<sub>2</sub> has the advantage of precise thickness control, film uniformity over large areas and conformal coating of complex 3-D surfaces [15]. Our initial demonstration had shown a marginal improvement in the slope of the transfer curves of GaN HEMTs with a slope of  $\sim 59$  mV/decade compared to the Boltzmann limit of  $\sim 66$  mV/decade at  $\sim 60$  °C, and suffered high gate leakage current under forward bias because of the use of Schottky gates. In addition, we had obtained sub-Boltzmann switching only in one scan direction of the gate bias [14]. In this paper, we overcome all these limitations by using AlGaIn/GaN MOS-HEMTs with ultralow leakage fabricated on Si platform, loaded with ALD VO<sub>2</sub> resistor at the

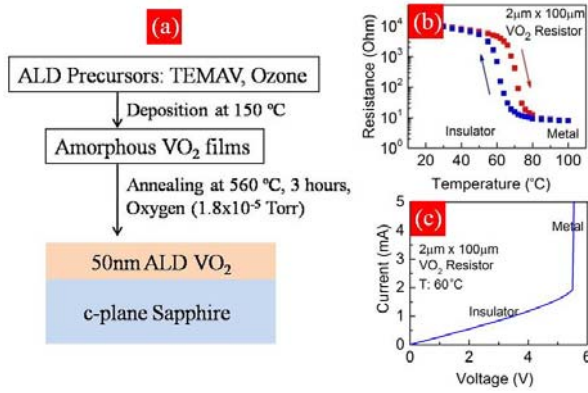


Fig. 2. (a) ALD process for VO<sub>2</sub>. (b) Temperature-driven MIT of the ALD-grown VO<sub>2</sub> showing a >3-order change in resistance at ~67 °C with the direction of hysteresis indicated. (c) Applied bias-driven MIT in 2- $\mu\text{m}$ -long, 100- $\mu\text{m}$ -wide VO<sub>2</sub> resistor measured at 60 °C, showing a current threshold of ~2 mA for the insulator–metal phase transition.

source (Fig. 1). Because of the low leakage current, we are able to demonstrate steep switching in both sweep directions with slopes of ~9.9 and ~28.2 mV/decade at ~60 °C, far below the Boltzmann limit (~66 mV/decade), using ALD VO<sub>2</sub>.

In Section II, we present the growth, fabrication, and characterization details of VO<sub>2</sub> resistors and AlGaIn/GaN MOS-HEMTs. The device characteristics of the control MOS-HEMTs and the phase-FETs are presented and discussed in a comparative fashion in Section III. In Section IV, we present the summary of this paper and offer our perspectives on the prospects of this technology for new functionalities in GaN-based low-power and high-power electronics, RF electronics, and photonics.

## II. ALD VO<sub>2</sub> RESISTORS AND NITRIDE MOS-HEMTS

VO<sub>2</sub> thin films used in this paper were grown by ALD using Tetrakis-ethylmethylamino-vanadium (TEMAV) and ozone as the precursors, as shown in Fig. 2(a). Films of 50-nm thickness were grown on c-plane sapphire substrates. The as-deposited films were amorphous, but were crystallized upon annealing. More details about ALD VO<sub>2</sub> growth process and optimization can be found in [15] and [16]. Optical lithography was used to pattern the VO<sub>2</sub> resistors. Ti/Au ohmic contacts were deposited using e-beam evaporation followed by mesa-isolation using SF<sub>6</sub>-based dry etching. VO<sub>2</sub> resistors of 100- $\mu\text{m}$  width and lengths varying from 2–20  $\mu\text{m}$  were fabricated. The measured temperature-driven MIT in the 2  $\mu\text{m}$  × 100  $\mu\text{m}$  VO<sub>2</sub> resistor is shown in Fig. 2(b). More than three-order change in resistance was measured when the layer went through the phase transition. The transition temperature was quite close to the bulk VO<sub>2</sub> MIT temperature (~67 °C) [12]. The measured applied bias-driven MIT in the 2  $\mu\text{m}$  × 100  $\mu\text{m}$  VO<sub>2</sub> resistor at 60 °C is shown in Fig. 2(c). The transition happens at an applied electric field of ~27 kV/cm and a current threshold of ~20  $\mu\text{A}/\mu\text{m}$ . Though this transition can be achieved at room temperature, an elevated temperature of 60 °C was used to keep the transition voltage lower.

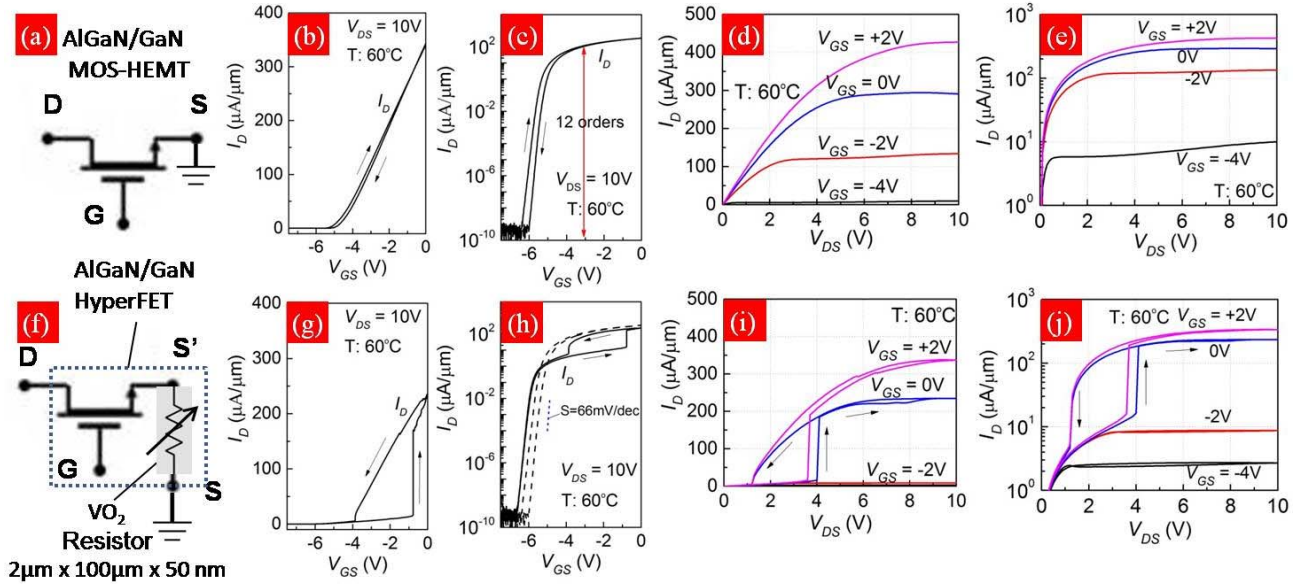
For the fabrication of the AlGaIn/GaN MOS-HEMTs, the following epilayer structure (Fig. 1) was used: 1-nm GaN cap/20-nm Al<sub>0.26</sub>Ga<sub>0.74</sub>N barrier/1-nm AlN spacer and a GaN channel on 1.3- $\mu\text{m}$  semiinsulating GaN buffer grown

on (111) Si substrates [17]. Ti/Al/Ni/Au (20/100/40/50 nm) metal stack annealed at 890 °C for 20 s in nitrogen ambient was used as the ohmic contact. The alloyed ohmic contacts showed low contact resistance of ~0.3  $\Omega \cdot \text{mm}$ . After mesa isolation (~300 nm), a gate dielectric stack of 5-nm SiN<sub>x</sub>/5-nm Al<sub>2</sub>O<sub>3</sub> was deposited using ALD followed by gate lithography and gate metal (40-nm Ni/100-nm Au) deposition. After the MOS-HEMT fabrication, a 2-D electron gas density of ~6.2 × 10<sup>12</sup> cm<sup>-2</sup> and electron Hall mobility ~1760 cm<sup>2</sup>/V · s were measured at room temperature. The circuit symbol for MOS-HEMT device denoting the source (S), gate (G) and drain (D) terminals is shown in Fig. 3(a). The measured  $I_D$ – $V_{GS}$  transfer characteristics and  $I_D$ – $V_{DS}$  characteristics of this MOS-HEMT device at 60 °C are shown in Fig. 3(b)–(e). The pinch-off voltage in Fig. 3(b) is measured to be –5 V, and ON/OFF ratio of more than 12 orders was achieved as shown in Fig. 3(c). Gate leakage current was below the measurement limit. The subthreshold slope of the MOS-HEMTs at low current levels is ~90 mV/decade, not far from the Boltzmann limit of 66 mV/decade at 60 °C. Fig. 3(d) shows good current saturation in MOS-HEMTs and current densities of ~0.4 mA/ $\mu\text{m}$ .

## III. NITRIDE PHASE-FETS

To realize the GaN phase-FET, a 200- $\mu\text{m}$ -wide MOS-HEMT with 3- $\mu\text{m}$  gate length, 2- $\mu\text{m}$  gate–source separation, and 4.5- $\mu\text{m}$  gate–drain separation were used. The source of the GaN MOS-HEMT was loaded with the 2  $\mu\text{m}$  × 100  $\mu\text{m}$  VO<sub>2</sub> resistor, as shown in Fig. 1 and Fig. 3(f). The measured transfer characteristics of the phase-FET are shown in Fig. 3(g) and (h), and the  $I_D$ – $V_{DS}$  characteristics are shown in Fig. 3(i). As mentioned earlier, all measurements were taken at 60 °C to keep the voltage threshold for VO<sub>2</sub> MIT low. Steep switching corresponding to the VO<sub>2</sub> insulator–metal transition (IMT) in the forward  $V_{GS}$  scan and the MIT in the reverse  $V_{GS}$  scan is clearly observed. Hysteresis is also observed in the transfer characteristics, as well as in the  $I_D$ – $V_{DS}$  characteristics of the phase-FET because of the VO<sub>2</sub> phase transition.

The drain current is suppressed in the phase-FETs compared to MOS-HEMTs at low gate and drain bias. This can be clearly observed on comparing log( $I_D$ ) versus  $V_{DS}$  curves for MOS-HEMT and phase-FET devices, as shown in Fig. 3(e) and (j). This behavior is expected because of the additional source resistance. When VO<sub>2</sub> source resistor is in the insulating state, effective  $V_{GS}$  bias of MOS-HEMT ( $V_{GS}'$ ) is reduced, leading to suppression of drain current. Even when the VO<sub>2</sub> resistor is in the metallic state at high drain/gate bias, it adds a small parasitic source resistance, reducing the intrinsic  $V_{GS}'$  and  $V_{DS}'$  bias, leading to reduction in the drain current. For example, at  $V_{GS} = +2$  V and  $V_{DS} = 10$  V, drain current in the MOS-HEMT is 427  $\mu\text{A}/\mu\text{m}$  while in the phase-FET; the drain current value under these bias conditions decreases to 337  $\mu\text{A}/\mu\text{m}$ . These results are consistent with the first demonstration of phase-FETs [10]. The apparent increase in the drain current of the phase-FET in [10] compared to the bare FET is because the results are plotted assuming iso-OFF condition (matched OFF-state current between bare FET and

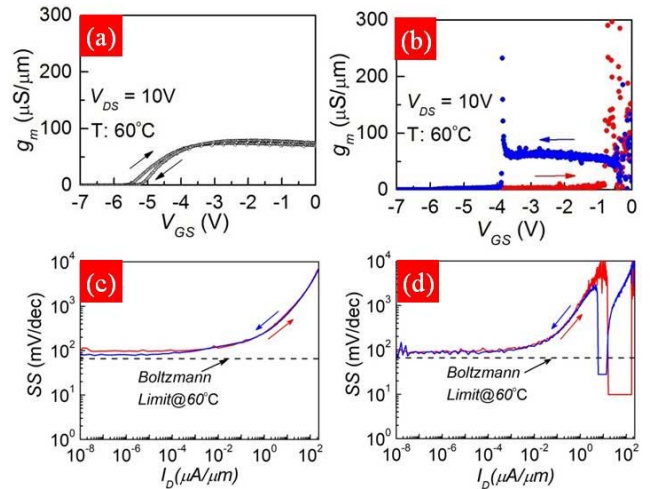


**Fig. 3.** (a) Circuit symbol for the stand-alone AlGaIn/GaN MOS-HEMT device. The source is grounded, and control voltages are applied to the gate and drain contacts. (b) Linear  $I_D$ - $V_{GS}$  transfer characteristics showing a threshold of  $\sim -5$  V. (c) Logarithmic  $I_D$ - $V_{GS}$  transfer characteristics showing 12-order on/off ratio. (d) and (e) Output  $I_D$ - $V_{DS}$  characteristics of the MOS-HEMT device at  $60^\circ\text{C}$  showing saturation. (f) Schematic circuit diagram of the  $\text{VO}_2$  resistor-loaded MOS-HEMT device or the GaN phase-FET. (g) Linear  $I_D$ - $V_{GS}$  transfer characteristics show a steep switching accompanied by hysteresis. (h) Logarithmic  $I_D$ - $V_{GS}$  transfer characteristics show the appearance of the sub-Boltzmann switching regimes, and the ultralow off-state leakage of the GaN phase-FET. The transfer characteristics of the control MOS-HEMT structure is shown in dashed lines for comparison. (i) and (j) Output  $I_D$ - $V_{DS}$  characteristics of the phase-FET device at  $60^\circ\text{C}$ , showing that there is a threshold drain voltage for the drain current turn on with steep characteristics.

the phase-FET). The FETs used in [10] have a small ON/OFF ratio of  $\sim 2$ – $3$  orders of magnitude, thus adding  $\text{VO}_2$  resistor in source lowers OFF current, and artificially gives enhanced drain currents when assuming iso-OFF condition (matched OFF-state current). The GaN MOS-HEMTs used in this paper have a large ON/OFF ratio of more than  $\sim 12$  orders and extremely low OFF currents (less than  $\sim 10^{-10} \mu\text{A}/\mu\text{m}$ ), thus adding  $\text{VO}_2$  resistor in source has no effect on the OFF current.

In Fig. 3(i), beyond  $V_{DS} = 4$  V, it can be observed that drain currents do not perfectly overlap in the forward and reverse sweeps even though  $\text{VO}_2$  has transitioned into the metallic state. This behavior arises because  $\text{VO}_2$  IMT/MIT is a first-order transition involving nucleation and growth of the opposite phase. During IMT transition metallic (rutile phase) domains nucleate and grow in size inside the insulating (monoclinic) phase matrix [18]–[20]. Once a conducting path is formed for current flow,  $\text{VO}_2$  resistance falls sharply. At this stage, some insulating domains can still be present in  $\text{VO}_2$ , which gradually convert to metallic phase with increasing bias. During the reverse sweep, most of the sample is metallic, leading to lower resistance/higher current (compared to forward voltage scan) and gradually insulating domains nucleate and grow to cause the MIT. This phenomena lead to nonoverlapping forward and reverse sweep currents.

In the subthreshold region of the MOS-HEMT, small hysteresis can be observed depending on the  $V_{GS}$  sweep direction [Fig. 3(c)]. This hysteresis is not present in the phase-FET [Fig. 3(h)] which is likely because the gate oxide is stressed differently in the phase-FET during forward and reverse scans compared to the MOS-HEMT. Due to large voltage drop



**Fig. 4.** Measured transconductance of (a) GaN MOS-HEMT and (b) GaN phase-FET as a function of gate bias. The GaN phase-FET has a much flatter  $g_m$  with hysteresis and peaks. The measured SS as a function of the drain current for (c) GaN MOS-HEMT and (d) GaN phase-FET device at  $60^\circ\text{C}$ . The GaN MOS-HEMT has an SS larger than the Boltzmann limit. The GaN phase-FET shows sub-Boltzmann switching at high currents with  $I_D$  well in excess of  $1 \mu\text{A}/\mu\text{m}$ .

across the  $\text{VO}_2$  resistor in the insulating state (forward scan), the intrinsic  $V_{GS}$ 's is more negative compared to the intrinsic  $V_{GS}$ 's in the reverse scan when  $\text{VO}_2$  is in the metallic state. This asymmetric gate oxide stress conditions can lead to different hysteresis behaviors in the phase-FET compared to the MOS-HEMT. Detailed investigation is needed for better understanding of this behavior.

Fig. 4(a) and (b) shows the measured transconductance ( $g_m$ ) of both the GaN MOS-HEMT and the GaN phase-FET. The phase-FET shows a nearly square-shaped  $g_m$  with hysteretic behavior, and sharp peaks at the sub-Boltzmann gate voltage points. Square-shaped  $g_m$  is obtained in the phase-FET for positive to negative  $V_{GS}$  sweep as VO<sub>2</sub> MIT truncates the typical nonlinear part of the  $g_m$  curve near pinch off. The switching slopes (SSs) for the GaN MOS-HEMT and the phase-FET are shown in Fig. 4(c) and (d), respectively. While the MOS-HEMT SS remains above the Boltzmann limit, very steep slopes of  $\sim 9.9$  and  $\sim 29.2$  mV/decade are measured in the phase-FET in forward and reverse  $V_{GS}$  scans, respectively. These values are significantly lower than the Boltzmann limit of 66 mV/decade at 60 °C. SS values for the MOS-HEMT at similar drain current levels are  $\sim 720$  mV/decade ( $I_D \sim 10 \mu\text{A}/\mu\text{m}$ ) and  $\sim 3200$  mV/dec ( $I_D \sim 100 \mu\text{A}/\mu\text{m}$ ). The phase-FET sub-Boltzmann SS values represent a decrease of  $\sim 25\times$  at  $I_D = 10 \mu\text{A}/\mu\text{m}$  and a decrease of more than  $\sim 300\times$  at  $I_D = 100 \mu\text{A}/\mu\text{m}$ . This experimental result is an initial proof of concept to access sub-Boltzmann limit modulation using ALD VO<sub>2</sub>. To move the steep transition gate voltages to the subthreshold regime for low-power digital switching, and potentially for memory-logic hybrids, it will be necessary to match the device geometries and the VO<sub>2</sub> impedance [21]. This will also enable shaping of the hysteresis. The GaN phase-FET exhibits a record high  $\sim 10^{12}$  ON/OFF ratio, which is attractive for low OFF-state leakage.

#### IV. CONCLUSION

In conclusion, we have demonstrated steep switching in ultralow leakage GaN phase-FETs realized by combining AlGaIn/GaN MOS-HEMTs on Si with ALD VO<sub>2</sub>. On-wafer integration of ALD VO<sub>2</sub> with GaN FETs [22], reducing hysteresis in phase-FETs, reducing VO<sub>2</sub> phase transition voltages at room temperature, and obtaining steep switching over larger drain current range will be the next steps toward realizing low-power GaN digital systems. VO<sub>2</sub> phase transition voltages can be reduced by scaling the lengths of VO<sub>2</sub> resistors to sub-micrometer regime [14], using strained VO<sub>2</sub> [23], [24] or by doping VO<sub>2</sub> [25], [26]. Several GaN device technologies beyond low-power digital can also benefit significantly from the technology demonstrated here. For example, low insertion loss, wideband VO<sub>2</sub> RF switches can directly be integrated with microwave transmission lines based on a GaN RF circuit platform [27], [28]. Integration with nitride LEDs and lasers can offer exciting potential for photonic triggering, reconfigurable mirrors, and plasmonics [29]–[31]. Finally, it can also enable new functionalities in GaN power electronics by offering new ways to realize lower ON-resistances and enhanced ON currents [10], [11].

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